

Fig. 3 Small-scale perturbation structure.

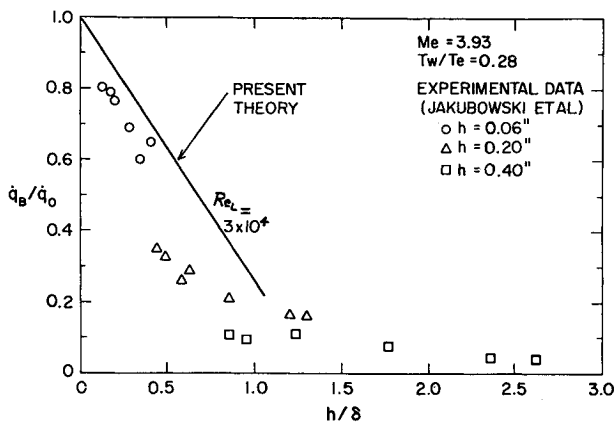


Fig. 5 Base heat transfer vs step height/boundary-layer thickness ratio.

The corresponding small-scale wall pressure solution has the form

$$p_w'/\gamma P_0 \approx -(4h/\pi y_s) M_0^2 (y^*) e^{kx/y_s} \quad (2)$$

in terms of the boundary-layer sonic height y_s where $k = 1$ and -3 for $x < 0$ and $x > 0$, respectively. This variation is continuous past the step (albeit with a discontinuous axial gradient) and acts to locally smooth out the large scale pressure jump illustrated in Fig. 1. Moreover, whereas the lateral pressure gradient in the large scale is negligibly small, it is quite large and rapidly varying (in fact comparable to the local axial pressure gradient) near the corner, as shown in Fig. 3.

Assuming laminar boundary-layer flow, the following useful closed-form approximation for the base pressure $p_B = p_w'(x \rightarrow 0^+)$ can be derived² from Eq. (2):

$$\frac{p_B}{p_0} \approx 1 - C_1 \frac{M_e^3 Re_L^{1/4} h}{(M_e^2 - 1)^{1/4} L} = 1 - C_1 \Lambda \quad (3)$$

where the constant C_1 depends on T_w/T_e . Equation (3) indicates that the quantity Λ is a basic correlating parameter for the base pressure. This is verified in Fig. 4, where several sets of experimental data^{4,5} are shown to correlate in terms of a single curve function of Λ over a moderately wide range of laminar flow conditions.

Solution of the energy equation in the heat-conducting disturbance sublayer near the wall gives the corresponding heat transfer behavior.² In the small-scale approximation near the step, this yields the following approximate laminar base heat-transfer expression:

$$\frac{q_B}{q_{w_0}} = 1 - C_2 \frac{M_e^2 Re_L^{5/8} h}{(M_e^2 - 1)^{1/2} L} \quad (4)$$

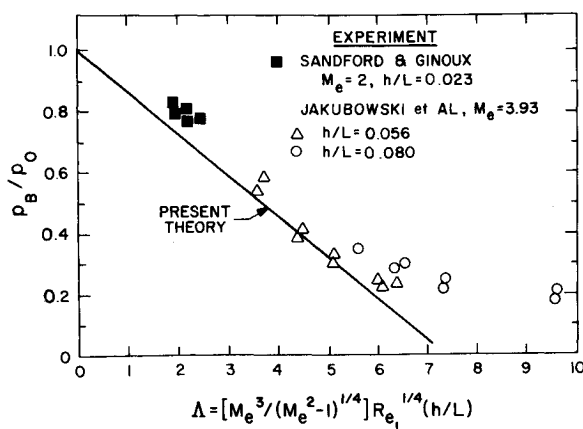


Fig. 4 Base pressure correlation for laminar supersonic flow.

which, as expected, exhibits a much stronger dependence on the Reynolds number than base pressure. This theoretical result is plotted vs h/δ in Fig. 5 along with some recent experimental measurements.⁴ Although the predicted linear behavior at small h/δ is confirmed, the theory overestimates the heat transfer noticeably more than it does base pressure. This is to be expected since the presence of a recirculation zone downstream of the step (which is not accounted for in the theory) significantly reduces base heating for the cited step heights.

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Generalized Sturm-Liouville Procedure for Composite Domain Anisotropic Transient Conduction Problems

JOSEPH PADOVAN*

University of Akron, Akron, Ohio

Introduction

THE transient and steady-state temperature distribution in composite configurations consisting of several distinct thermally anisotropic subdomains have numerous applications to heat-transfer problems in re-entry vehicles, air frames, nuclear reactors, and the like. Apart from the purely numerical approaches such as the finite element¹⁻³ and difference procedures,⁴ several purely analytical techniques⁴⁻⁸ are also available. A comprehensive survey of the above noted analytical procedures has been reported by Ozisik.⁴ With few exceptions,^{9,10} most applications of analytical techniques have been limited to isotropic 1-D laminated domain problems. This is partly due to the added analytical complexity caused by anisotropy as well as the general lack of attention given to such thermal material properties. Due to the increased usage of inherently thermally anisotropic materials in a variety of applications, this situation is changing.

In this context, the purpose of the present Note will be to extend the relatively simple and straightforward Vodicka-Tittle⁴⁻⁶ orthogonal expansion technique to 3-D configurations consisting of finitely many distinct fully anisotropic subdomains. As will be seen in the development, the procedure developed herein is based on a 3-D piecewise weighted orthogonality

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* Associate Professor of Mechanical Engineering. Member AIAA.